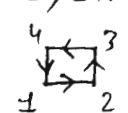


Thermodynamics Properties

- | | |
|--|---|
| <p style="text-align: center;"><u>Intensive</u></p> <ul style="list-style-type: none"> * Not depend on mass or quantity of system → Density, Volume → Temp., Pressure → Heat, Refractive Index → Surface tension | <p style="text-align: center;"><u>Extensive</u></p> <ul style="list-style-type: none"> * depend on quantity of system Mass, Volume, energy, enthalpy, entropy, etc. |
|--|---|
- * Some Intensive property derive from extensive property
 e.g. specific Volume, specific Heat, density

- * If X and Y intensive then $X+Y$, XY , $\frac{X}{Y}$, $\frac{\partial X}{\partial Y}$ will also intensive
- * If X and Y extensive then $X+Y$ will be extensive
 But $\frac{X}{Y}$ and $\frac{\partial X}{\partial Y}$ are Intensive

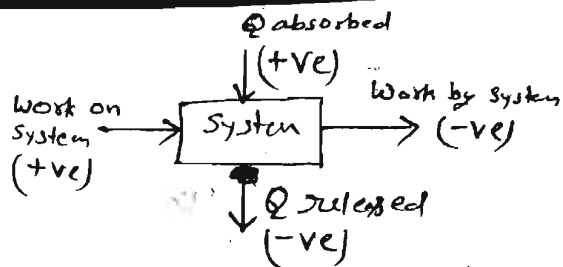
Thermodynamic Process

- | | |
|---------------------|--|
| ① Isothermal $dT=0$ | ② Isobaric process $dP=0$ |
| ③ Isochoric $dV=0$ | ④ Adiabatic $dQ=0$
(or isentropic process)
(isentropic) |
- ④ Cyclic $dU=0$, $dH=0$
- 
1 → 2 → 3 → 4
↑
All point function changes are zero

* Reversible process is ideal not in actual, it take infinite time for completion, maximum work obtained in this type process.

⇒ First Law of Thermodynamic

$$\Delta U = Q + W$$

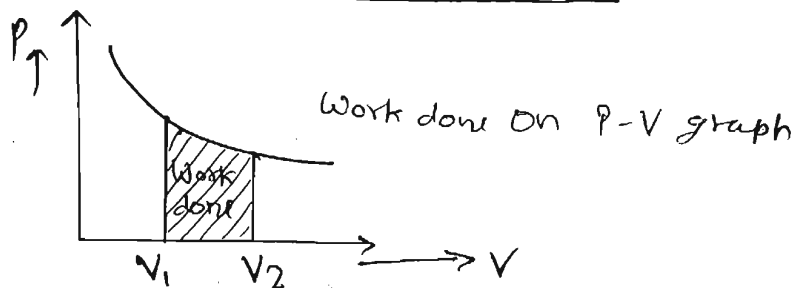


Sign conversion

- ① for work done on the system W take $+ve$
- ② for work done by system W take $-ve$
- ③ Heat absorbed by system Q take $+ve$
- ④ Heat released by system Q take $-ve$

⇒ Work done If gas enclosed in a volume change its from V_1 to V_2 , work done.

$$W = \int_{V_1}^{V_2} P dV$$



⇒ Point function Property Value depend on Initial & final stage. Internal energy, enthalpy, entropy

⇒ Path function change in property value depend on path e.g., work, heat.

⇒ Thermodynamics criteria for irreversible process

- $(\Delta G)_{T,P} < 0$
- $(\Delta A)_{T,V} < 0$
- $(\Delta U)_{S,V} < 0$
- $(\Delta H)_{S,P} < 0$

⇒ Limitation of First Law

- ① 1st Law does not tell the direction of Heat flow
 "It says that the heat extracted from a cold object by further reducing its temp. But it does not say that we actually need an external machine for doing so"
- ② It does not deny that all Heat converted into work. While any machine does not give 100% efficiency
- ③ It does not tell the difference b/w spontaneous and Non spontaneous process and why Natural spontaneous process are irreversible.
- ④ It does not tell why certain chemical reactions do not proceed to completion.

⇒ Heat Capacity Relation

$$\text{Heat capacity} = \frac{\Delta Q}{\Delta T}$$

Required energy for 1°C change in temp

at const Volume

$$C_v = \frac{Q_v}{\Delta T}$$

$$\Delta U = Q + W = Q + \int P dV$$

$$\Delta U = Q \quad \text{so} \quad C_v = \frac{\Delta U}{\Delta T}$$

at const pressure

$$C_p = \frac{Q_p}{\Delta T}, \quad H = U + PV$$

$$C_p = \frac{\Delta H}{\Delta T}$$

$$\Delta U = Q + W = Q + \int_{V_1}^{V_2} P dV$$

$$\Delta U = Q + P(V_2 - V_1)$$

$$U_2 - U_1 = Q + P V_2 - P V_1$$

$$(U_2 + P V_2) - (U_1 + P V_1) = Q$$

$$H_2 - H_1 = Q$$

$$\Delta H = Q_p \quad \text{at const } P$$

$$W = -\int P dV \quad \text{in expansion}$$

then

$$C_v = \left(\frac{\partial U}{\partial T} \right)_V$$

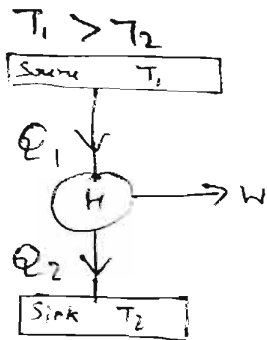
$$C_p = \left(\frac{\partial H}{\partial T} \right)_P$$

$$\gamma = \frac{C_p}{C_v}$$

$$C_p - C_v = R$$

⇒ Heat Engine Performance of heat engine measured by its efficiency

$$\eta = \frac{\text{Work output}}{\text{Heat Input}} = \frac{W}{Q_1} = \frac{Q_1 - Q_2}{Q_1} = 1 - \frac{Q_2}{Q_1}$$



$$\eta = 1 - \frac{Q_2}{Q_1}$$

Since $Q \propto T$

$$\eta = 1 - \frac{T_2}{T_1}$$

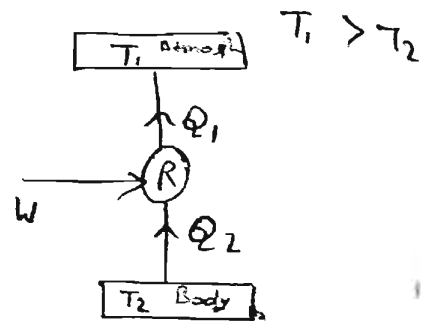
⇒ Refrigerator

$$(\text{COP})_R = \frac{\text{Heat removed}}{\text{Work Input}} = \frac{Q_2}{W}$$

Now energy balance

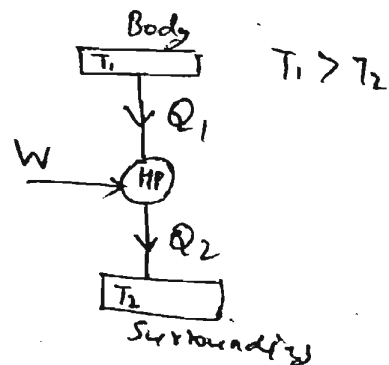
$$W + Q_2 = Q_1 \Rightarrow W = Q_1 - Q_2$$

$$(\text{COP})_R = \frac{Q_2}{Q_1 - Q_2} = \frac{T_2}{T_1 - T_2}$$



⇒ Heat pump

$$(\text{COP})_{HP} = \frac{Q_1}{W} = \frac{Q_1}{Q_1 - Q_2} = \frac{T_1}{T_1 - T_2}$$



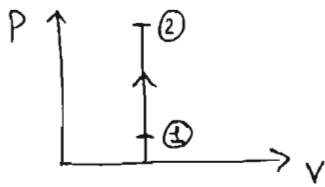
$$HP = R + 1$$

⇒ Isochoric Process ($dV = 0$)

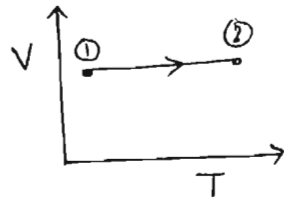
$PV = nRT$
 $P \propto T$

* Work done $W = \int P dV = 0$

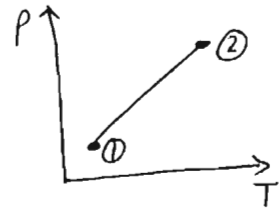
* Graph of Isochoric process



P-V graph



V-T graph



* Internal energy change $\Delta U = Q + W = Q_V + 0$

$\Delta U = n C_V \Delta T$

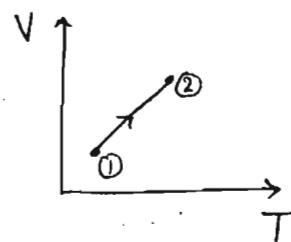
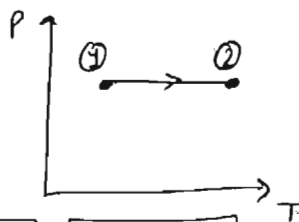
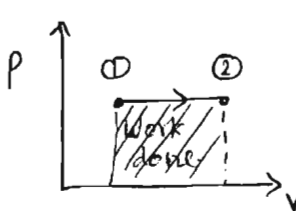
⇒ Isoobaric Process ($dP = 0$)

* Work done $W = \int P dV = P \int dV = P \Delta V$ where $\Delta V = V_2 - V_1$

$W = nR \Delta T$

$\Delta T = T_2 - T_1$

* Graph of Isoobaric process



* Internal energy $\Delta U = n C_V \Delta T$ $\Delta H = n C_P \Delta T$

⇒ Isothermal Process ($dT = 0$)

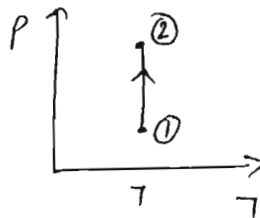
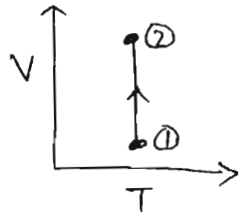
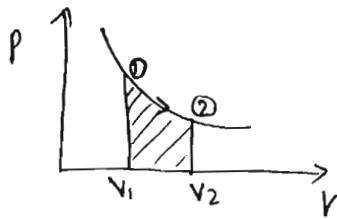
* $\Delta U = n C_V \Delta T = 0$, $\Delta H = 0$

* $\Delta U = Q + W$ But $\Delta U = 0$ then $Q = -W$

* $W = -nRT \ln\left(\frac{V_2}{V_1}\right) = RT \ln\left(\frac{P_2}{P_1}\right)$ Where $V_2 > V_1$

* $W = nRT \ln\left(\frac{V_1}{V_2}\right) = nRT \ln\left(\frac{P_2}{P_1}\right)$ Where $V_1 > V_2$

* Curve of Isothermal Process



Adiabatic Process ($dQ = 0$) (Reversible process) (expansion)

* $\Delta U = Q + W = W$

* Work done $W = \Delta U = n c_v \Delta T$

But $W = -P \Delta V = -P dV$

then $c_v dT = -P dV = -\frac{RT dV}{V}$ (from $PV = RT$)

$$\int_1^2 \frac{dT}{T} = -\frac{R}{c_v} \int_1^2 \frac{dV}{V}$$

$$\ln \frac{T_2}{T_1} = -\frac{R}{c_v} \ln \frac{V_2}{V_1} = \frac{R}{c_v} \ln \frac{V_1}{V_2}$$

But $c_p - c_v = R \Rightarrow \frac{c_p}{c_v} - 1 = \frac{R}{c_v} \Rightarrow (\gamma - 1) = \frac{R}{c_v}$

then $\ln \frac{T_2}{T_1} = (\gamma - 1) \ln \frac{V_1}{V_2}$

$$\left(\frac{T_2}{T_1} \right) = \left(\frac{V_1}{V_2} \right)^{\gamma - 1} \Rightarrow T_2 V_2^{\gamma - 1} = T_1 V_1^{\gamma - 1}$$

and also $PV^\gamma = \text{const}$

$$W = c_v \Delta T = \frac{R}{\gamma - 1} \Delta T = \frac{R(T_2 - T_1)}{\gamma - 1} = \frac{P_2 V_2 - P_1 V_1}{\gamma - 1}$$

\Rightarrow Irreversible, Adiabatic, expansion. workdone

$\Delta U = W = 0$ when external pressure zero

$W = -P_{\text{ext}} (V_2 - V_1)$ when, $P_{\text{external}} < P_{\text{gas}}$

Important Points (All for reversible, closed, Ideal gas system)

① Adiabatic workdone = $c_v \Delta T = \frac{R \Delta T}{\gamma - 1} = \frac{R(T_2 - T_1)}{\gamma - 1} = \frac{P_2 V_2 - P_1 V_1}{\gamma - 1}$

② $\Delta U = W$ in Adiabatic, $P(V)^\gamma = \text{const}$ $\gamma = \frac{c_p}{c_v}$

$c_p - c_v = R$

③ In Isochoric process workdone = 0

④ In Isothermal $W = nRT \ln \frac{V_1}{V_2} = nRT \ln \frac{P_2}{P_1}$

⑤ In Throttling process these condition apply

a) Adiabatic $P(V)^\gamma = \text{const}$

b) $\Delta H = 0$, $W = 0$, $Q = 0$

c) Isenthalpic, Irreversible

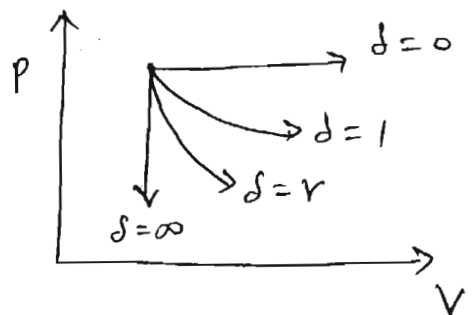
⑥ Polytropic process $P(V)^\delta = \text{const}$

i) when $\delta = 0 \Rightarrow$ Isobaric

ii) when $\delta = 1 \Rightarrow$ Isothermal

iii) when $\delta = \gamma \Rightarrow$ Adiabatic

iv) when $\delta = \infty \Rightarrow$ Isochoric

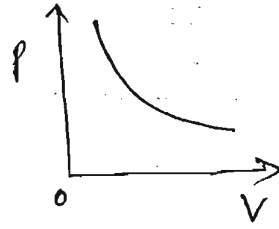


Workdone $W = \frac{P_2 V_2 - P_1 V_1}{\delta - 1}$

Heat $Q = \left(\frac{\delta - \gamma}{\delta - 1} \right) n c_v (T_2 - T_1)$

⇒ Workdone in Polytropic

$$P(V)^\delta = \text{const} = k$$



$$W = - \int P dV$$

$$= - \int_1^2 \frac{k}{V^\delta} dV = - \int_1^2 k V^{-\delta} dV$$

$$= -k \left[\frac{V^{-\delta+1}}{-\delta+1} \right]_{V_1}^{V_2}$$

$$= -k \left(\frac{V_2^{1-\delta}}{1-\delta} \right)_{V_1} = \frac{-k}{1-\delta} (V_2^{1-\delta} - V_1^{1-\delta})$$

$$= \frac{-k}{1-\delta} \left(\frac{V_2}{V_2^\delta} - \frac{V_1}{V_1^\delta} \right)$$

$$= \frac{-k}{1-\delta} \left(\frac{V_2 P_2}{k} - \frac{V_1 P_1}{k} \right) = \frac{-1}{1-\delta} (P_2 V_2 - P_1 V_1)$$

$$W = - \left(\frac{P_2 V_2 - P_1 V_1}{1-\delta} \right)$$

where $\delta \neq 1$

Its not apply for Isothermal

⇒ Q (Heat) $\Delta U = Q + W$

$$Q = \Delta U - W$$

$$W = \left(\frac{P_2 V_2 - P_1 V_1}{\delta - 1} \right) = \frac{n R (T_2 - T_1)}{\delta - 1}$$

then $Q = n C_V \Delta T - n R \frac{(T_2 - T_1)}{\delta - 1}$

$$= n C_V (T_2 - T_1) - n C_V \frac{(\gamma - 1)(T_2 - T_1)}{\delta - 1}$$

$$C_p - C_v = R$$

$$\frac{C_p}{C_v} - 1 = \frac{R}{C_v}$$

$$(\gamma - 1) = \frac{R}{C_v}$$

$$C_v = \frac{R}{(\gamma - 1)}$$

$$Q = n C_V (T_2 - T_1) \left[1 - \frac{\gamma - 1}{\delta - 1} \right]$$

$$Q = \left(\frac{\delta - \gamma}{\delta - 1} \right) n C_V (T_2 - T_1)$$

Joule Thomson effect

* Joule Thomson coeff.

$$\mu = \left(\frac{\partial T}{\partial P} \right)_H = \frac{1}{C_p} \left(\frac{2a}{RT} - b \right)$$

* This process is isenthalpic process.

* the temp at which μ change sign its known as Inversion temp. of the gas.

* $\mu > 0$ Gas will show cooling effect in expansion
 $\mu < 0$ Gas will show warming effect in expansion

* For an ideal gas $\mu = 0$

* for real gas μ take any value including zero
($-\infty$ to $+\infty$)

* At inversion temp $\mu = 0$ then $b = \frac{2a}{RT_i}$

OR
$$T_i = \frac{2a}{Rb}$$

Entropy (unit J/K-mol)

- * Entropy is state function, Extensive property.
- * Entropy is randomness or disorder of molecules in system
- * Difficult to determine the actual entropy of system becauz randomness is not a fixed parameter.
- * Entropy change at absolute temp (T)

$$\Delta S = \int \frac{dq_{rev}}{T}$$

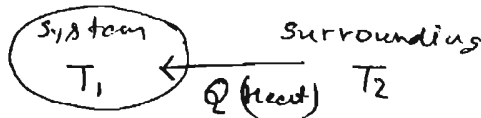
Heat absorbed $Q \rightarrow +ve$ then $\Delta S = +ve$
 Heat released $Q \rightarrow -ve$ then $\Delta S = -ve$

- * By \uparrow se temp \rightarrow Randomness \uparrow se \rightarrow Entropy \uparrow se
- * In solid molecules are fixed, so no randomness and entropy will be zero at $(T=0K)$ IIIrd Law
- * $\Delta S_{total} > 0$ for spontaneous, $\Delta S_{total} < 0$ for Nonspontaneous

⊛ IInd Law $\Delta S_{total} \geq 0$ | V. Jump

⊛ In Irreversible

$$\Delta S > 0$$



$$\begin{aligned} \Delta S_{total} &= \Delta S_{system} + \Delta S_{surroundings} \\ &= +\frac{Q}{T_1} - \frac{Q}{T_2} = Q \left(\frac{1}{T_1} - \frac{1}{T_2} \right) \end{aligned}$$

$$\Delta S > 0$$

⊛ Reversible

$$\Delta S_{total} = 0$$

$$\Delta S_{total} = \Delta S_{system} + \Delta S_{surroundings}$$

$$= +\frac{Q}{T} - \frac{Q}{T} = 0$$



* When entropy is function of T, V

$$\Delta S = n C_V \ln\left(\frac{T_2}{T_1}\right) + nR \ln\left(\frac{V_2}{V_1}\right)$$

①

* When entropy is function of T, P

$$\Delta S = n C_P \ln\left(\frac{T_2}{T_1}\right) + nR \ln\left(\frac{P_1}{P_2}\right)$$

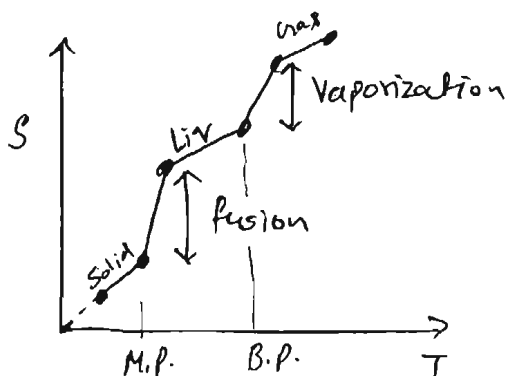
②

* Entropy $\left\{ \begin{array}{l} \uparrow \text{sc in expansion} \\ \downarrow \text{sc in compression} \end{array} \right\}$

* Entropy change of Mix of Gas

$$\Delta S_{\text{mix}} = -\sum R n_i \ln n_i$$
$$\Delta S_{\text{mix}}^{\text{mole}} = -\sum R n_i \ln n_i$$

\Rightarrow Tephigraph (Variation of Entropy vs Temp.)



* $S \uparrow$ sc with T

* Entropy sudden \uparrow sc during phase change

Moller diagram is H vs S

PVT Relation

- ⊗ Ideal gas Law $PV = nRT$
- ⊗ at low pressure, high temp gas is considered Ideal
- ⊗ Vanderwaals eqⁿ $\left(P + \frac{a}{V^2}\right)(V-b) = RT$
- ⊗ At critical condition ($T = T_c$, $P = P_c$)
 $\left(\frac{\partial P}{\partial V}\right)_T = \left(\frac{\partial^2 P}{\partial V^2}\right)_T = 0$
- ⊗ Compressibility factor
 $Z = \frac{PV}{RT}$
- ⊗ Virial eqⁿ $Z = 1 + \frac{B}{V} + \frac{C}{V^2}$
- ⊗ Another form of virial eqⁿ
 $Z = 1 + B'P + C'P^2$

⇒ First Law $\Delta U = Q + W$

⇒ Work done in reversible process $W = -\int P dV$

⇒ Work done in irreversible process $W = -P_r \Delta V$, ($\Delta V = V_2 - V_1$)

⇒ Heat capacities $C = \frac{dQ}{dT}$

⇒ Heat capacities at const. volume $C_v = \left(\frac{\partial U}{\partial T}\right)_V$ where $dQ = dU$

⇒ Heat capacities at const. pressure. $C_p = \left(\frac{\partial H}{\partial T}\right)_P$ where $dQ = dH$

Thermodynamics Relation

⇒ Gibbs free energy $G = H - TS$

⇒ Helmholtz free energy $A = U - TS$

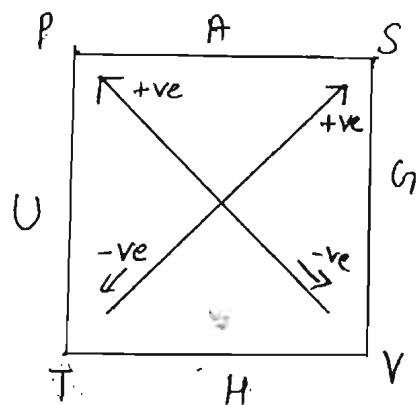
⇒ Fundamental Property Relation

$$dU = +T dS - P dV$$

$$dH = +T dS + V dP$$

$$dG = +V dP - S dT$$

$$dA = -S dT - P dV$$



⇒ Short Method (U H G A)

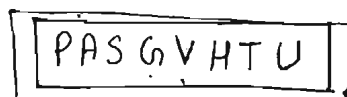
$dU = dW + dQ$ — 1st Law.

$$dU = -P dV + T dS$$

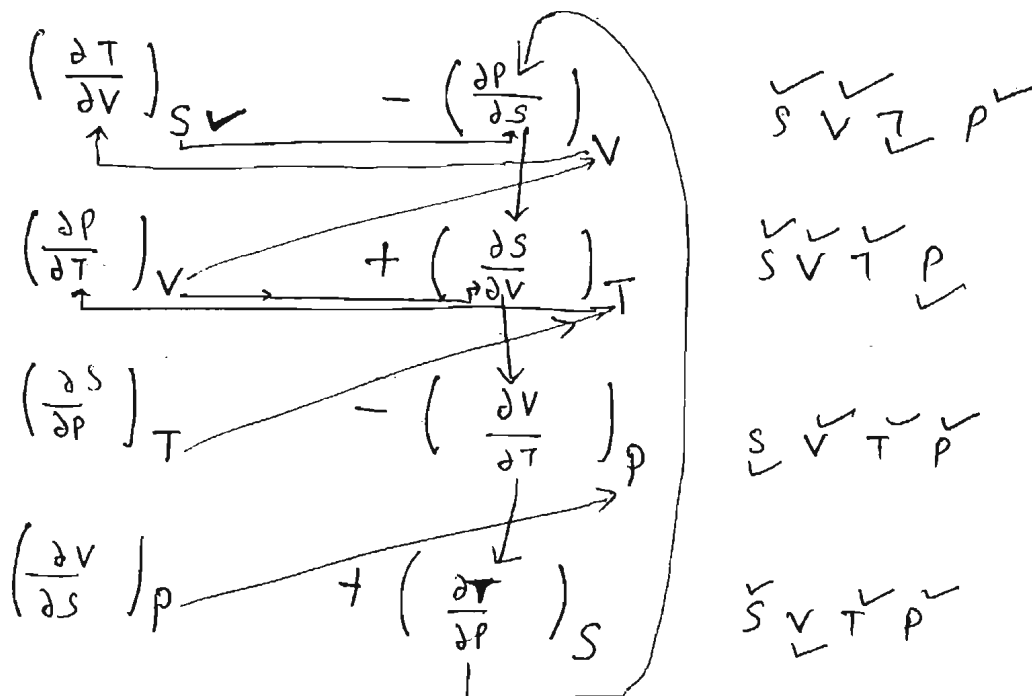
$$dH = T dS + V dP$$

$$dG = V dP - S dT$$

$$dA = -S dT - P dV$$



⇒ Short trick for write Maxwell eqⁿ (S V T P)



(T V S P) Maxwell eqnⁿ

Best Method

$$\begin{aligned} \left(\frac{\partial S}{\partial P}\right)_T &= - \left(\frac{\partial V}{\partial T}\right)_P & \left(\frac{\partial S}{\partial P}\right)_T &= - \left(\frac{\partial V}{\partial T}\right)_P \\ \left(\frac{\partial P}{\partial T}\right)_V &= + \left(\frac{\partial S}{\partial V}\right)_T & \left(\frac{\partial P}{\partial T}\right)_V &= + \left(\frac{\partial S}{\partial V}\right)_T \\ \left(\frac{\partial T}{\partial V}\right)_S &= - \left(\frac{\partial P}{\partial S}\right)_V & \left(\frac{\partial T}{\partial V}\right)_S &= - \left(\frac{\partial P}{\partial S}\right)_V \\ \left(\frac{\partial V}{\partial S}\right)_P &= + \left(\frac{\partial T}{\partial P}\right)_S & \left(\frac{\partial V}{\partial S}\right)_P &= + \left(\frac{\partial T}{\partial P}\right)_S \end{aligned}$$

⇒ For Real gas Workdone in Adiabatic expansion

$$W = nRT \ln \left(\frac{V_1 - nb}{V_2 - nb} \right) + an^2 \left(\frac{1}{V_1} - \frac{1}{V_2} \right)$$

* Change in Internal energy

$$\Delta U = an^2 \left(\frac{1}{V_1} - \frac{1}{V_2} \right)$$

* Heat

$$Q = nRT \ln \left(\frac{V_2 - nb}{V_1 - nb} \right)$$

⇒ For Isobaric, Isochoric workdone is same for Ideal gas.

$$\Rightarrow \text{Mach Number} = \frac{\text{Velocity of fluid}}{\text{Velocity of sound in fluid}}$$

* For a closed system

$$d(nG) = (nV)dP - (nS)dT$$

* For open system of constant or variable composition

$$d(nG) = (nV)dP - (nS)dT + \sum \mu_i dn_i$$

where $\mu_i = \left(\frac{\partial (nG)}{\partial n_i} \right)_{P, T, n_j}$

* during Phase change $P, T,$ and n are constant.

* Clapeyron Equation
(where two phase are in equilibrium then it apply)

$$\frac{dP^*}{dT} = \frac{\Delta H^{lv}}{T \Delta V^{lv}} \quad \frac{dP}{dT} = \frac{\Delta H}{T(V_2 - V_1)}$$

* Clausius - Clapeyron Equation (approximate on low pressure) Valid for

Assumption $\Delta V^{lv} = V^v = \frac{RT}{P^*}$

$$\frac{dP^*/P^*}{dT/T^2} = \frac{\Delta H^{lv}}{R} \quad \text{i.e.} \quad \ln \left(\frac{P_2^*}{P_1^*} \right) = \frac{\Delta H}{R} \left(\frac{1}{T_1} - \frac{1}{T_2} \right)$$

* Antoine Equation

$$\ln P^* = A - \frac{B}{T+C}$$

* Mollier diagram H-S graph of Water

* Residual property for gas

* Excess property for liq

* Isothermal compressibility

$$\alpha = -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_T$$

* coefficient of volume expansion

$$\beta = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_P$$

$$\frac{dV}{V} = \beta dT - \alpha dP$$

Residual Property

* Residual Property (V^R, U^R, H^R, S^R, G^R) for gases only
extensive

* $M^R = M - M^{ig}$

↓

actual value of Thermodynamics
Property of Substance at
given (T, P)

→

Value of Thermodynamic property of
an Ideal gas at
same T, P

for example (Volume)

$$V^R = V - V^{ig} = V - \frac{RT}{P} \quad \left(PV = RT \right) \text{ Ideal}$$

But for real gas

$$V = \frac{zRT}{P} \text{ eqn of state}$$

$$V^R = \frac{zRT}{P} - \frac{RT}{P} \Rightarrow V^R = \frac{RT}{P}(z-1)$$

* $\frac{G^R}{RT} = \int_0^P (z-1) \frac{dP}{P}$ ✓

* $\frac{G^R}{RT} = z - 1 - \ln z + \frac{1}{RT} \int_{\infty}^V \left(\frac{RT}{v} - P \right) dv$ for real gas ✓

* $U^R = \int_{\infty}^V \left[T \left(\frac{\partial P}{\partial T} \right)_v - P \right] dv$ ✓ All relation from $PV = zRT$

* $H^R = \int_0^P \left[v - T \left(\frac{\partial v}{\partial T} \right)_P \right] dP + RT(z-1)$ ✓ eqn of state

* $S^R = \int_0^P \left[\frac{R}{P} - \left(\frac{\partial v}{\partial T} \right)_P \right] dP$ ✓

Residual Property from Virial Eqⁿ

$$* \frac{H^R}{RT_c} = P_r \left[\left(0.083 - \frac{1.097}{T_r^{1.6}} \right) + \omega \left(0.139 - \frac{0.894}{T_r^{4.2}} \right) \right]$$

$$* \frac{S^R}{R} = -P_r \left[\frac{0.675}{T_r^{2.6}} + \omega \left(\frac{0.722}{T_r^{5.2}} \right) \right]$$

Residual Property from Cubic Eqⁿ of state
Vander Waals

$$* \boxed{U^R = -\frac{a}{V}}$$

$$\boxed{S^R = R \ln \frac{P(V-b)}{RT}}$$

$$* \boxed{H^R = \frac{bRT}{V-b} - \frac{2a}{V}}$$

⇒ Partial Molar properties

$$\bar{M}_i \equiv \left[\frac{\partial(nM)}{\partial n_i} \right]_{P, T, n_j}$$

\bar{M}_i = Partial properties

M = solution properties

M_i = Pure species properties

⇒ Relation b/w Molar and partial Molar property

$$nM = f(T, P, n_1, n_2, \dots, n_i)$$

$$d(nM) = \left(\frac{\partial(nM)}{\partial P} \right)_T dP + \left(\frac{\partial(nM)}{\partial T} \right)_P dT + \sum_j \left[\frac{\partial(nM)}{\partial n_j} \right]_{P, T, n_j} dn_j$$

at const T, P and Binary Solⁿ

$$d(nM) = n dM + M dn \quad \text{and} \quad n_i = x_i n$$

$$dn_i = x_i dn + n dx_i$$

then

$$n dM + M dn = n \left(\frac{\partial M}{\partial P} \right)_{T, n} dP + n \left(\frac{\partial M}{\partial T} \right)_{P, n} dT + \sum \bar{M}_i (x_i dn + n dx_i)$$

$$\left[dM - \left(\frac{\partial M}{\partial P} \right)_{T, n} dP - \left(\frac{\partial M}{\partial T} \right)_{P, n} dT - \sum \bar{M}_i dx_i \right] n + \left[M - \sum x_i \bar{M}_i \right] dn = 0$$

\downarrow
 \downarrow

Zero
Zero

$$\boxed{dM = \left(\frac{\partial M}{\partial P} \right)_{T, n} dP + \left(\frac{\partial M}{\partial T} \right)_{P, n} dT + \sum \bar{M}_i dx_i} \quad \text{--- (1)}$$

$$\boxed{M = \sum x_i \bar{M}_i} \Rightarrow \boxed{dM = \sum x_i d\bar{M}_i + \sum \bar{M}_i dx_i} \quad \text{--- (2)}$$

from ① & ②

$$\boxed{\left(\frac{\partial M}{\partial P} \right)_{T, n} dP + \left(\frac{\partial M}{\partial T} \right)_{P, n} dT - \sum x_i d\bar{M}_i = 0}$$

at const T, P $\boxed{\sum x_i d\bar{M}_i = 0}$ V. Prop

$$\Rightarrow \bar{M}_i = \left[\frac{\partial(MM)}{\partial n_i} \right]_{P, T, n_j}$$

which yields partial properties from total properties

$$\Rightarrow M = \sum n_i \bar{M}_i$$

which yields total properties from partial properties

\Rightarrow Method to find partial properties

$$M = n_1 \bar{M}_1 + n_2 \bar{M}_2 \quad \text{--- A}$$

$$dM = \underline{n_1 d\bar{M}_1} + \bar{M}_1 dn_1 + \underline{n_2 d\bar{M}_2} + \bar{M}_2 dn_2$$

But $\sum n_i d\bar{M}_i = 0$ at const. T, P

$$\rightarrow \boxed{dM = \bar{M}_1 dn_1 + \bar{M}_2 dn_2} \quad \text{and} \quad \boxed{n_1 d\bar{M}_1 + n_2 d\bar{M}_2 = 0}$$

$$n_1 + n_2 = 1 \Rightarrow \boxed{dn_1 = -dn_2}$$

$$dM = -\bar{M}_1 dn_2 + \bar{M}_2 dn_2 = dn_2 (\bar{M}_2 - \bar{M}_1)$$

$$\frac{dM}{dn_2} = \bar{M}_2 - \bar{M}_1 \quad \text{--- (B)}$$

Now eqⁿ (A) $M = (1-n_2)\bar{M}_1 + n_2\bar{M}_2$
 $= \bar{M}_1 + \bar{M}_2(n_2) - \bar{M}_1(n_2)$
 $M = \bar{M}_1 + n_2(\bar{M}_2 - \bar{M}_1)$

$$\boxed{M = \bar{M}_1 + n_2 \frac{dM}{dn_2}}$$

$$M = n_1 \bar{M}_1 + (1-n_1)\bar{M}_2 = n_1(\bar{M}_1 - \bar{M}_2) + \bar{M}_2$$

$$M = -n_1 \frac{dM}{dn_2} + \bar{M}_2 = \bar{M}_2 + n_1 \frac{dM}{dn_1}$$

$$\boxed{dn_2 = -dn_1}$$

$$\boxed{M = \bar{M}_2 + n_1 \frac{dM}{dn_1}}$$

$$\& \quad \boxed{M = \bar{M}_1 + n_2 \frac{dM}{dn_2}}$$

* In Ideal Gas Mixture

$$\textcircled{1} \quad \bar{H}_i^{ig} = H_i^{ig}$$

$$\textcircled{2} \quad \bar{S}_i^{ig} = S_i^{ig} - R \ln y_i \quad \text{and} \quad \bar{G}_i^{ig} = \bar{H}_i^{ig} - T \bar{S}_i^{ig}$$

$$\textcircled{3} \quad \bar{G}_i^{ig} = G_i^{ig} + RT \ln y_i \quad \leftarrow$$

$$\textcircled{4} \quad H^{ig} = \sum y_i H_i^{ig}$$

$$\textcircled{5} \quad S^{ig} = \sum y_i S_i^{ig} - R \sum y_i \ln y_i$$

$$\textcircled{6} \quad G^{ig} = \sum y_i G_i^{ig} + RT \sum y_i \ln y_i$$

$$\textcircled{7} \quad G_i^{ig} = f_i(T) + RT \ln P$$

$$\textcircled{8} \quad \bar{G}_i^{ig} = f_i(T) + RT \ln (y_i P) = \mu_i$$

(Chemical Potential)

Chemical Potential (Intensive Property)

$$\mu_i = \left(\frac{\partial G}{\partial n_i} \right)_{P, T, n_j}$$

Chemical Potential is partial derivative of H, A, G of j th component with respect to composition

⑨ Effect of P on chemical potential

$$\bar{V}_i = \left(\frac{\partial \mu_i}{\partial P} \right)_{T, n_j}$$

⑩ Effect of temp on μ

$$-\bar{S}_i = \left(\frac{\partial \mu_i}{\partial T} \right)_{P, n_j}$$

Ideal gas Mixture Model

$$\Delta H = \sum n_i (\bar{H}_i - H_i^\circ)$$

① Entropy change of mixing

$$\Delta S_{mix} = -nR \sum x_i \ln x_i$$

② Gibbs free energy change of mixing

$$\Delta G_{mix} = nRT \sum x_i \ln x_i$$

③ Enthalpy change of mixing for Ideal gas Mix.

$$\Delta H_{mix} = 0$$

④ Volume change of mixing for Ideal gas Mix

$$\Delta V_{mix} = 0$$

⑤ Molar free energy of gas

$$\mu_i = \mu_i^\circ + RT \ln x_i$$

⑥ $\Delta U_{mix}^{ideal} = 0$

Fugacity

is fictitious pressure which is used for real gas

We know that for ideal gas, closed system

$$dG^{ig} = VdP - SdT$$

at const. temp. $\left(\frac{\partial G^{ig}}{\partial P}\right)_T = V$

$$\left(dG^{ig}\right)_T = VdP = \frac{nRT}{P}dP = nRT d(\ln P) \quad \text{--- (1)}$$

Integrating this eqn

$$G = G^{\circ} + nRT \ln P \quad \text{--- (2)}$$

But its apply on real gas so eqn (1) introduced new variable for real gas

$$dG = nRT d(\ln f) \quad f = \text{fugacity} \quad \text{--- (3)}$$

from eqn ~~(2)~~ (2) - (1)

$$d(G - G^{ig}) = nRT d(\ln f) - nRT d(\ln P)$$

$$\boxed{dG^R = RT d \ln \left(\frac{f}{P}\right)}$$

$G^R = G - G^{ig}$
Residual

Fugacity
coeff: ϕ

$$\phi = \frac{f}{P}$$

$$dG^R = RT d(\ln \phi)$$

On Integration $G - G^{ig} = G^R = RT \ln \phi + C$

Put $\phi = 1$, $G^R = 0$ then $C = 0$ so

$$G^R = RT \ln \phi \Rightarrow$$

$$\boxed{\frac{G^R}{RT} = \ln \phi}$$

$$\boxed{\frac{G^R}{RT} = \ln \phi = \sum n_i \ln \bar{\phi}_i}$$

⇒ Fugacity Coeff from Compressibility factor

$$z = \frac{PV}{RT} = 1 + \frac{BP}{RT}$$

$$z - 1 = \frac{BP}{RT}$$

$$\ln \phi = \int_0^P \frac{z-1}{P} dP$$

⇒ from virial eqⁿ of state

$$\ln \phi = \ln \frac{f}{P} = \int_0^{P_r} (z-1) \frac{dP}{P_r} = (B^0 + w B^1) \frac{P_r}{T_r}$$

⇒ Gibbs Duhem eqⁿ

$$\textcircled{1} \quad \sum n_i d\bar{M}_i = 0$$

$$\textcircled{2} \quad \sum n_i d\mu_i = 0$$

$$\textcircled{3} \quad \sum n_i d\ln a_i = 0$$

$$\textcircled{4} \quad \sum n_i d(\ln f_i) = 0$$

$$\textcircled{5} \quad \sum n_i d(\ln a_i) = 0$$

$$\textcircled{6} \quad \sum n_i d(\ln f_i) = 0$$

⇒ Fugacity of gas Mix

$$\ln \phi = \sum n_i \ln \bar{\phi}_i$$

⇒ +ve deviation $\gamma_i > 1$

-ve deviation $\gamma_i < 1$

No deviation $\gamma_i = 1$

Excess property of Mixture

$$\left[\begin{array}{l} \text{Actual property of} \\ \text{sol}^n \text{ at given} \\ \text{at } T, P \end{array} \right] - \left[\begin{array}{l} \text{Property it would have} \\ \text{as an ideal gas} \\ \text{at same } T, P \end{array} \right] = \text{Excess property} \\ \text{of sol}^n \text{ at given } T, P$$

$$M - M^{id} = M^E$$

⇒ Excess free Energy

$$\frac{G^E}{RT} = \sum n_i \ln \gamma_i = \sum n_i \frac{\bar{G}_i^E}{RT}$$

$$\ln \gamma_i = \frac{\bar{G}_i^E}{RT} = \left[\frac{\partial (n \frac{G^E}{RT})}{\partial n_i} \right]_{T, P, n_j}$$

V. Imp.

⇒ criteria for Equilibrium

- ① Isolated system $dS_{U,V} \leq 0$ (4)
- ② at const T, V $dA \leq 0$
- ③ at const T, P $dG_{T,P} \leq 0$

⇒ criteria for Phase

- ① Chemical potential in each phase equal
- ② fugacity in each phase equal of all component of Mix.
- ③ $dn = 0$

⇒ Phase rule for Non reacting system

① Phase Rule

$$F = C - P + 2$$

where F = no. of Independent intensive variable

- * In this phase rule only intensive variable include.
- * '2' is show that Temp, Pressure → Intensive.

$$* \text{ DOF} = \text{No. of Variable} - \text{No. of Eq}^n$$

- * If F = 1 that is one independent intensive property is required to be specified
- * at Triple point F = 0 It implies that none of the properties of a pure substance at triple point can be varied.

② DUHEM'S Theorem

In this extensive + intensive both consider.

no. of Intensive variable = $P(C-1) + 2$

no. of Extensive " = $P \left[\begin{matrix} \text{mass in} \\ \text{each phase} \end{matrix} \right]$

total Variable = $P(C-1) + 2 + P$
 $= 2 + CP$

no. of Intensive eqⁿ relation = $(P-1)C$

no. of extensive " " = C

total " " = $(P-1)C + C$

$$\text{DOF} = [2 + CP] - [(P-1)C + C]$$

$$= 2 + CP - [PC + C + C]$$

$$= 2$$

Acco. to this only two^v Variable are required to define state

③ Phase rule derive

consider system (closed) with P = Phase, C = Component

In closed system required only (C-1) component composition, other one is remainder. so

no. of Composition Variable in 'P' Phase = $P(C-1)$

two variable is T & P

$$\text{total Variable} = P(C-1) + 2$$

In P-phase, C-component only (P-1)C no. of equilibrium phase eqⁿ so

$$\text{no. of Eq}^n = (P-1)C$$

$$F = N_v - N_E = P(C-1) + 2 - (P-1)C$$

$$= PC - P + 2 - PC + C$$

$$F = C - P + 2$$

at least one variable extensive ← F = 1

both must be ← F = 0
 Extensive

VLE Point

⇒ If ① = Higher Volatile, ② = less Volatile than all ways

$$y_1 > x_1 \text{ and } y_2 < x_2$$

⇒ Raoult's Law

* for liq phase composition

$$\boxed{\sum y_i = 1} \quad \boxed{P = x_1 P_1^* + x_2 P_2^*} \quad \text{and} \quad \boxed{y_i P = x_i P_i^*}$$

* for Vapor Phase System

$$\boxed{\frac{1}{P} = \frac{y_1}{P_1^*} + \frac{y_2}{P_2^*}} \quad \boxed{\sum x_i = 1}$$

⇒ Modified Raoult's Law

Vapor = Ideal

$$\boxed{y_i P = x_i \gamma_i P_i^*}$$

γ_i = Activity coeff

Liq = Real

Vapor = Real

$$\boxed{y_i \phi_i P = x_i \gamma_i P_i^*}$$

ϕ_i = Fugacity coeff.

Liq = Real

⇒ Deviation from Raoult's Law

+ve deviation

-ve deviation

① $P_{total} > P_{Raoult \text{ Law}}$

① $P_{total} < P_{Raoult \text{ law}}$

② $p_A > p_A^* x_A$

② $p_A < p_A^* x_A$

③ $A \text{---} B < A \text{---} A, B \text{---} B$
(Adhesive force) Cohesive force

③ $A \text{---} B > A \text{---} A, B \text{---} B$
Adhesive Cohesive force

④ $\Delta H > 0$ Endothermic process

④ $\Delta H < 0$ Exothermic process

⑤ $\Delta S > 0, \Delta V > 0$

⑤ $\Delta S < 0, \Delta V < 0$

example

MeOH + H₂O
EtOH + H₂O
Acetone + CS₂

example

HCl + H₂O
HNO₃ + H₂O
CHCl₃ + H₂O
Acetone + CHCl₃

Azetrope [Constant Boiling Mixture]

⇒ Each Azetrope has a characteristic boiling point.

* B.P. of Azetrope > Boiling point of its every constituents. for (-ve) Azetrope
Maximum boiling Mix

* B.P. of Azetrope < B.P. of its every constituents for (+ve) Azetrope
Minimum boiling Mix.

⇒ Maximum boiling Mixture

⇒ Minimum boiling Mixture

① $(B.P.)_{AZ} > (B.P.)_{\text{Pure component}}$

$(B.P.)_{AZ} < (B.P.)_{\text{Pure component}}$

② B.P. $\propto \frac{1}{V.P.}$ So $(V.P.)_{AZ} < (V.P.)_{\text{Pure component}}$

$(V.P.)_{AZ} > (V.P.)_{\text{Pure component}}$

③ Large -ve deviation

Large +ve deviation

④ -ve Azetrope

+ve Azetrope

⑤ $\ln \gamma_i < 0, \gamma_i < 1$

$\ln \gamma_i > 0, \gamma_i > 1$

⑥ e.g. (HCl + water)
 20.2% 79.8%

e.g. Pyridine + water
 59% 41%

⇒ Margules eqnⁿ (3-suffix)

$$\ln \gamma_1 = \gamma_2^2 [A + 2(B-A)\gamma_1]$$

$$\ln \gamma_2 = \gamma_1^2 [B + 2(A-B)\gamma_2]$$

⇒ Margules eqnⁿ (2-suffix)

$$\ln \gamma_1 = A\gamma_2^2, \quad \ln \gamma_2 = A\gamma_1^2$$

⇒ VLE at High pressure (K-value)

$$K = \frac{y_i}{x_i} = \frac{P_i^*}{P}$$

for Modified Law
 Raoult's

$$K = \frac{y_i P_i^*}{P}$$

⇒ For Binary Mixture, the ratio two k-values of two component is called relative volatility

$$\alpha = \frac{k_i}{k_j} = \frac{\frac{y_i}{n_i}}{\frac{y_j}{n_j}}$$

⇒ At 100% Relative Humidity [dew point = current temp]

⇒ Dew point is also known saturation point

⇒ When dew point temp falls below freezing it is often called frost point

Flash Calculation

Consider a binary Mixture, feed rate F mol/h

total Material balance

$$F = V + L$$

Material Balance for component

$$Fz_i = Vy_i + Ln_i$$

Put $y_i = k_i n_i$ then

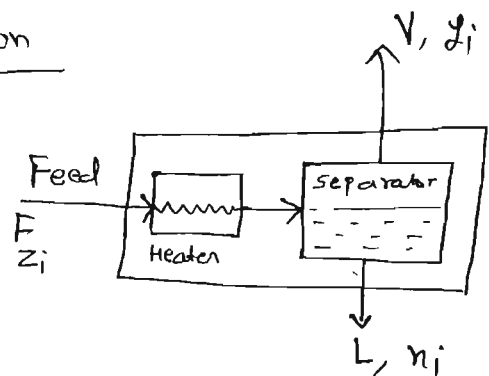
$$Fz_i = V k_i n_i + L n_i$$

$$= n_i (V k_i + L)$$

$$n_i = \frac{Fz_i}{L + V k_i}$$

$\sum n_i = 0$ then

$$\sum \frac{Fz_i}{L + V k_i} = 1$$



Put $n_i = \frac{y_i}{k_i}$ then

$$Fz_i = Vy_i + \frac{L y_i}{k_i} = y_i \left[V + \frac{L}{k_i} \right]$$

$$y_i = \frac{Fz_i}{\left[V + \frac{L}{k_i} \right]} \quad \sum y_i = 1$$

then

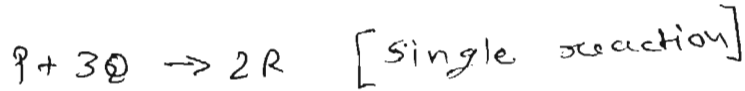
$$\sum \frac{Fz_i}{\left[V + \frac{L}{k_i} \right]} = 1$$

Chemical reaction Equilibria

⇒ RXN Coordinate (ϵ)

* It represent progress of chemical reaction along a path way

* This is one dimensional coordinate



$$V_R = 2, V_P = -1, V_Q = -3$$

[product +ve
Reactant -ve
Inert zero]

$$\frac{dn_i}{V_i} = d\epsilon \Rightarrow \int_{n_{i0}}^{n_i} dn_i = \int_0^\epsilon V_i d\epsilon$$

$$n_i - n_{i0} = V_i \epsilon \Rightarrow \boxed{n_i = n_{i0} + V_i \epsilon}$$

total mole $n = \sum n_i = \sum n_{i0} + \epsilon \sum V_i$

$$\boxed{n = n_0 + \epsilon V}$$

where $\boxed{V = \sum V_i}$

$$\boxed{n_0 = \sum n_{i0}}$$

$$\boxed{y_i = \frac{n_i}{n} = \frac{n_{i0} + V_i \epsilon}{n_0 + \epsilon V}}$$

* In Multi reaction

$$y_i = \frac{n_i}{n} = \frac{n_{i0} + \sum_j V_{i,j} \epsilon_j}{n_0 + \sum_j V_j \epsilon_j}$$

$i = \text{component}$

$j = \text{reaction}$

① Equilibrium criterion of a chemical rxn

$$(dG)_{T,P} \leq 0$$

② At equilibrium

$$(dG)_{T,P} = 0$$

③ Equilibrium constant K_c is depend on only temp.

④ $\alpha A + \beta B \rightleftharpoons \sigma S + \tau T$

① $K_c = \frac{[C_S]^\sigma [C_T]^\tau}{[C_A]^\alpha [C_B]^\beta}$ (in terms of concⁿ)

② $K_a = \frac{a_S^\sigma a_T^\tau}{a_A^\alpha a_B^\beta} = \prod (a_i)^{\nu_i}$ (in terms of activity)

③ $K_f = \prod \left(\frac{f_i}{f_i^\circ} \right)^{\nu_i}$ (in terms of fugacity)

④ $K_p = \prod (P_i)^{\nu_i}$ (in terms of partial Pr)

⑤ $K_p = K_c (RT)^{\Delta n}$ (where $\Delta n = \text{mol of product} - \text{mol of reactant}$)

⑥ If $A + B \rightleftharpoons C$ equilibrium const = K_c

then $C \rightleftharpoons A + B$ equilibrium const = $\frac{1}{K_c}$

$2A + 2B \rightleftharpoons 2C$ eqⁿ const = $(K_c)^2$

① Equilibrium constant K_c in terms of concn

② Equilibrium constant K_p in terms of partial pressure

③ $K_p = K_c (RT)^{\Delta n}$

④ $K_a = \prod (a_i)^{v_i} = \prod \left(\frac{f_i}{f_i^\circ} \right)^{v_i} = \prod (f_i)^{v_i} = K_f$
 $= \prod (\bar{p}_i)^{v_i} = K_p$

$K_a = K_f = K_p$
for ideal system $K_c = K_a$ becauz activity coeff = 1

⇒ Homogeneous gas phase rxn ($p^\circ = 1$)

$K_a = K_p = K_f = K_y p^{\sum v_i}$ (K_p and K_c both find)

and $K_y = \prod (y_i)^{v_i}$

$K_p = K_c (RT)^{\Delta n}$

⇒ Homogeneous liq phase
(only K_c find)

$K_a = K_n = \prod (n_i)^{v_i}$

⊗ Activity of solid component = one

⇒ Heterogeneous phase (find K_p) (only gas phase term consider)

⇒ $\Delta G^\circ = -RT \ln K$

⇒ Van't Hoff eqn (effect of temp on K)

$\ln \frac{K_2}{K_1} = \frac{\Delta H^\circ}{R} \left(\frac{1}{T_1} - \frac{1}{T_2} \right)$

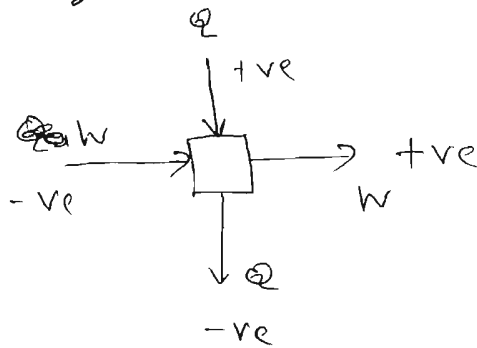
$\frac{\partial (\ln K)}{\partial T} = \frac{\Delta H^\circ}{RT^2}$

* Ientropic process means

- ① reversible
- ② Adiabatic
- ③ entropy constant
- ④ used for cooling
- ⑤ used for real + Ideal gas
- ⑥ $\mu_s = \left(\frac{\partial T}{\partial P}\right)_s$

* Ienthalpic process means (Joule Thomson effect)

- ① enthalpy const.
- ② Initial temp. should < Inversion temp for cooling
- ③ Adiabatic
- ④ used for real gas only
- ⑤ $\mu_{JT} = \left(\frac{\partial T}{\partial P}\right)_{dH}$



* Adiabatic

$$\begin{cases}
 PV^\gamma = \text{const.} \\
 TV^{\gamma-1} = \text{const.} \Rightarrow \frac{T}{P^{1-\gamma}} \\
 T^\gamma P^{1-\gamma} = \text{const.} \\
 T^\gamma \propto P^{1-\gamma}
 \end{cases}$$

then used 1st Law \Rightarrow

$$\Delta U = Q - W$$

$$W_{ad, rev} = \frac{R(T_2 - T_1)}{1 - \gamma} = C_V (T_1 - T_2)$$

$$\Delta U_{ad, rev} = C_V (T_2 - T_1)$$

$$dQ = 0$$

These are for
Rev and Adiabatic
Ientropic